Instructional Objectives:

At the end of this lesson, the students should have the knowledge of:

- Failure theories applied to thick walled pressure vessels.
- Variation of wall thickness with internal pressure based on different failure theories.
- Failure criterion of prestressed thick cylinders.
- Comparison of wall thickness variation with internal pressure for solid wall, single jacket and laminated thick walled cylinders.
- Failure criterion for thick walled cylinders with autofrettage.

9.3.1 Application of theories of failure for thick walled pressure vessels.

Having discussed the stresses in thick walled cylinders it is important to consider their failure criterion. The five failure theories will be considered in this regard and the variation of wall thickness to internal radius ratio $t/r_i$ or radius ratio $r_o/r_i$ with $p/\sigma_{yp}$ for different failure theories would be discussed. A number of cases such as $p_o=0$, $p_i=0$ or both non-zero $p_o$ and $p_i$ are possible but here only the cylinders with closed ends and subjected to an internal pressure only will be considered, for an example.

9.3.1.1 Maximum Principal Stress theory

According to this theory failure occurs when maximum principal stress exceeds the stress at the tensile yield point. The failure envelope according to this failure mode is shown in figure-9.3.1.1.1 and the failure criteria are given by $\sigma_1 = \sigma_2 = \pm \sigma_{yp}$. If $p_o=0$ the maximum values of circumferential and radial stresses are given by

$$\sigma_{\theta(\text{max})}\bigg|_{r=r_i} = p_i \left(\frac{r_o^2 + r_i^2}{r_o^2 - r_i^2}\right) \quad \sigma_{r(\text{max})}\bigg|_{r=r_i} = -p_i$$

(1)
Here both $\sigma_0$ and $\sigma_r$ are the principal stresses and $\sigma_0$ is larger. Thus the condition for failure is based on $\sigma_0$ and we have

$$
\frac{p_i}{r_i} \left( \frac{r_0^2}{r_0^2 - r_i^2} \right)^2 = \sigma_{yp}
$$

where $\sigma_{yp}$ is the yield stress.

This gives

$$
t = \frac{1}{r_i} \left[ \sqrt{\frac{1 + \frac{p_i}{\sigma_{yp}}}{1 - \frac{p_i}{\sigma_{yp}}} - 1} \right]
$$

(2)

9.3.1.1.1F- Failure envelope according to Maximum Principal Stress Theory.

9.3.1.2 Maximum Shear Stress theory

According to this theory failure occurs when maximum shear stress exceeds the maximum shear stress at the tensile yield point. The failure envelope according to this criterion is shown in figure 9.3.1.2.1 and the maximum shear stress is given by

$$
\tau_{max} = \frac{\sigma_1 - \sigma_2}{2}
$$

where the principal stresses $\sigma_1$ and $\sigma_2$ are given by
Here $\sigma_1$ is tensile and $\sigma_2$ is compressive in nature. $\tau_{\text{max}}$ may therefore be given by

$$\tau_{\text{max}} = p_i \frac{r_o^2}{r_o - r_i}$$

(3)

and since the failure criterion is $\tau_{\text{max}} = \sigma_{yp} / 2$ we may write

$$\frac{t}{r_i} = \sqrt{\frac{1}{1 - 2 \left( \frac{p_i}{\sigma_{yp}} \right)}} - 1$$

(4)

9.3.1.2.1F- Failure envelope according to Maximum Shear Stress theory.

9.3.1.3 Maximum Principal Strain theory

According to this theory failure occurs when the maximum principal strain exceeds the strain at the tensile yield point.
\[ \varepsilon_1 = \frac{1}{E} \left( \sigma_1 - \nu (\sigma_2 + \sigma_3) \right) = \varepsilon_{yp} \] and this gives \( \sigma_1 - \nu (\sigma_2 + \sigma_3) = \sigma_{yp} \)

where \( \varepsilon_{yp} \) and \( \sigma_{yp} \) are the yield strain and stress respectively. Following this the failure envelope is as shown in **figure-9.3.1.3.1**. Here the three principle stresses can be given as follows according to the standard 3D solutions:

\[ \sigma_1 = \sigma_0 = p_1 \frac{r_0 - r_i}{r_0 + r_i} , \quad \sigma_2 = \sigma_r = -p_1 \quad \text{and} \quad \sigma_3 = \sigma_z = \frac{p_1 r_i^2}{r_0 - r_i} \]

(5)

The failure criterion may now be written as

\[ p_1 \left( \frac{r_0 - r_i}{r_0 + r_i} + \nu - \frac{\nu r_i^2}{r_0 - r_i} \right) = \sigma_{yp} \] and this gives

\[ \frac{t}{r_i} = \sqrt{\frac{1 + (1 - 2\nu) p_i / \sigma_{yp}}{1 - (1 + \nu) p_i / \sigma_{yp}}} - 1 \]

(6)

**9.3.1.3.1F- Failure envelope according to Maximum Principal Strain theory**
9.3.1.4 Maximum Distortion Energy Theory

According to this theory if the maximum distortion energy exceeds the distortion energy at the tensile yield point failure occurs. The failure envelope is shown in figure-9.3.1.4.1 and the distortion energy $E_d$ is given by

$$ E_d = \frac{1 + \nu}{6E} \left\{ \left( \sigma_1 - \sigma_2 \right)^2 + \left( \sigma_2 - \sigma_3 \right)^2 + \left( \sigma_3 - \sigma_1 \right)^2 \right\} $$

Since at the uniaxial tensile yield point $\sigma_2 = \sigma_3 = 0$ and $\sigma_1 = \sigma_{yp}$

$E_d$ at the tensile yield point $= \frac{1 + \nu}{3E} \sigma_{yp}^2$

We consider $\sigma_1 = \sigma_0$, $\sigma_2 = \sigma_r$ and $\sigma_3 = \sigma_z$ and therefore

$$ \sigma_1 = p_i \frac{r_o^2 + r_i^2}{r_o - r_i} $$

$$ \sigma_r = -p_i, \quad \sigma_z = \frac{p_i r_i^2}{r_o - r_i} $$

(7)

The failure criterion therefore reduces to

$$ \frac{p_i}{\sigma_{yp}} = \frac{1}{\sqrt{3}} \left( \frac{r_o^2 - r_i^2}{r_o^2 / r_i} \right) $$

which gives

$$ \frac{t}{r_i} = \sqrt{\frac{1 - \sqrt{3} p_i / \sigma_{yp}}{1 - \sqrt{3} p_i / \sigma_{yp}}} - 1 $$

(8)
9.3.1.4.1F- Failure envelope according to Maximum Distortion Energy Theory

Plots of \( \pi/\sigma_{yp} \) and \( t/r_i \) for different failure criteria are shown in figure-9.3.1.4.2.

9.3.1.4.2F- Comparison of variation of \( \frac{\pi}{\sigma_{yp}} \) against \( \frac{t}{r_i} \) for different failure criterion.

The criteria developed and the plots apply to thick walled cylinders with internal pressure only but similar criteria for cylinders with external...
pressure only or in case where both internal and external pressures exist may be developed. However, on the basis of these results we note that the rate of increase in \( \frac{p_i}{\sigma_{yp}} \) is small at large values of \( t/r_i \) for all the failure modes considered. This means that at higher values of \( p_i \) small increase in pressure requires large increase in wall thickness. But since the stresses near the outer radius are small, material at the outer radius for very thick wall cylinders are ineffectively used. It is therefore necessary to select materials so that \( \frac{p_i}{\sigma_{yp}} \) is reasonably small. When this is not possible prestressed cylinders may be used.

All the above theories of failure are based on the prediction of the beginning of inelastic deformation and these are strictly applicable for ductile materials under static loading. Maximum principal stress theory is widely used for brittle materials which normally fail by brittle fracture.

In some applications of thick cylinders such as, gun barrels no inelastic deformation can be permitted for proper functioning and there design based on maximum shear stress theory or maximum distortion energy theory are acceptable. For some pressure vessels a satisfactory function is maintained until inelastic deformation that starts from the inner radius and spreads completely through the wall of the cylinder. Under such circumstances none of the failure theories would work satisfactorily and the procedure discussed in section lesson 9.2 is to be used.

### 9.3.1.5 Failure criteria of pre-stressed thick cylinders

Failure criteria based on the three methods of pre-stressing would now be discussed. The radial and circumferential stresses developed during shrinking a hollow cylinder over the main cylinder are shown in figure-9.3.1.5.1.
9.3.1.5.1F- Distribution of radial and circumferential stresses in a composite thick walled cylinder subjected to an internal pressure.

Following the analysis in section 9.2 the maximum initial (residual) circumferential stress at the inner radius of the cylinder due to the contact pressure $p_s$ is

$$\sigma_{o|t=r_i} = -2p_s \frac{r_s^2}{r_o - r_s}$$

and the maximum initial (residual) circumferential stress at the inner radius of the jacket due to contact pressure $p_s$ is

$$\sigma_{o|t=r_s} = p_s \frac{r_o^2 + r_s^2}{r_o - r_s}$$

Superposing the circumferential stresses due to $p_i$ (considering the composite cylinder as one) the total circumferential stresses at the inner radius of the cylinder and inner radius of the jacket are respectively
\[ \sigma_{\theta}^{\text{max}} = -2p_s \frac{r_s^2}{r_s - r_i^2} + p_i \frac{r_o^2 + r_i^2}{r_o - r_i^2} \]

\[ \sigma_{\theta}^{\text{max}} = p_s \frac{r_o^2 + r_s^2}{r_o - r_s^2} + p_i \frac{r_i^2}{r_s^2 + \left( \frac{r_o^2 + r_s^2}{r_s^2 - r_i^2} + \frac{r_i^2}{r_i^2 - r_i^2} \right)} \]

These maximum stresses should not exceed the yield stress and therefore we may write

\[ -2p_s \frac{r_s^2}{r_s - r_i^2} + p_i \frac{r_o^2 + r_i^2}{r_o - r_i^2} = \sigma_{yp} \] (9)

\[ p_s \frac{r_o^2 + r_s^2}{r_o - r_s^2} + p_i \frac{r_i^2}{r_s^2 + \left( \frac{r_o^2 + r_s^2}{r_s^2 - r_i^2} + \frac{r_i^2}{r_i^2 - r_i^2} \right)} = \sigma_{yp} \] (10)

It was shown in section-9.2 that the contact pressure \( p_s \) is given by

\[ p_s = \frac{E \delta}{\frac{r_s^2}{r_s^2 + \frac{r_o^2 + r_s^2}{r_o - r_s^2} + \frac{r_i^2}{r_i^2 - r_i^2}}} \] (11)

From (9), (10) and (11) it is possible to eliminate \( p_s \) and express \( t/r_i \) in terms of \( p_i/\sigma_{yp} \) and this is shown graphically in figure-9.3.1.5.2.
9.3.1.5.2F- Plot of $p_i/\sigma_{yp}$ vs $t/r_i$ for laminated multilayered, single jacket and solid wall cylinders.

This shows that even with a single jacket there is a considerable reduction in wall thickness and thus it contributes to an economic design.

As discussed earlier autofrettage causes yielding to start at the inner bore and with the increase in pressure it spreads outwards. If now the pressure is released the outer elastic layer exerts radial compressive pressure on the inner portion and this in turn causes radial compressive stress near the inner portion and tensile stress at the outer portion. For a given fluid pressure during autofrettage a given amount of inelastic deformation is produced and therefore in service the same fluid pressure may be used without causing any additional elastic deformation.

The self hooping effect reaches its maximum value when yielding just begins to spread to the outer wall. Under this condition the cylinder is said to have reached a fully plastic condition and the corresponding internal fluid pressure is known as fully plastic pressure, say, $p_f$. This pressure may be found by using the reduced equilibrium equation (3) in section- 9.2.1 which is reproduced here for convenience

\[
\sigma_\theta = \sigma_r + r \frac{d\sigma_r}{dr}
\]

(12)

Another equation may be obtained by considering that when the maximum shear stress at a point on the cylinder wall reaches shear yield value $\tau_{yp}$ it remains constant even after further yielding. This is given by

\[
\frac{1}{2} (\sigma_\theta - \sigma_r) = \tau_{yp}
\]

(13)

However experiments show that fully plastic pressure is reached before inelastic deformation has spread to every point on the wall. In fact Luder’s lines appear first. Luder’s lines are spiral bands across the cylinder wall such that the material between the bands retains elasticity. If the cylinder
is kept under fully plastic pressure for several hours uniform yielding across the cylinder wall would occur.

This gives \( \frac{d\sigma_r}{dr} = \frac{2\tau_{yp}}{r} \) and on integration we have

\[
\sigma_r = 2\tau_{yp} \log r + c
\]

Applying the boundary condition at \( r = r_o \) \( \sigma_r = 0 \) we have

\[
\sigma_r = 2\tau_{yp} \log \left( \frac{r}{r_o} \right) \quad \text{and} \quad \sigma_\theta = 2\tau_{yp} \left( 1 + \log \left( \frac{r}{r_o} \right) \right)
\]

(14)

Also applying the boundary condition at \( r = r_i \) \( \sigma_r = -p_f \) we have

\[
p_f = -2\tau_{yp} \log \left( \frac{r}{r_o} \right)
\]

(15)

Since the basic equations are independent of whether the cylinders are open or closed ends, the expressions for \( \sigma_r \) and \( \sigma_\theta \) apply to both the conditions. The stress distributions are shown in figure-9.3.1.5.3.

**9.3.1.5.3F- Stress distribution in a thick walled cylinder with autofrettage**

If we roughly assume that \( 2\tau_{yp} = \sigma_{yp} \) we have
\[
\frac{p_f}{\sigma_{yp}} = -\log \left( \frac{r_i}{r_o} \right)
\]

(16)

The results of maximum principal stress theory and maximum shear stress theory along with the fully plastic results are replotted in figure 9.3.1.5.4 where we may compare the relative merits of different failure criteria. It can be seen that cylinders with autofrettage may endure large internal pressure at relatively low wall thickness.

9.3.1.5.4F- Plots of \( \frac{p}{\sigma_{yp}} \) vs \( \frac{r_i}{r_o} \) for maximum shear stress theory, maximum principal stress theory and maximum autofrettage.
Finally it must be remembered that for true pressure vessel design it is essential to consult Boiler Codes for more complete information and guidelines. Pressure vessels can be extremely dangerous even at relatively low pressure and therefore the methodology stated here is a rough guide and should not be considered to be a complete design methodology.

9.3.2 Problems with Answers

Q.1: Determine the necessary thickness of the shell plates of 2.5m diameter boiler with the internal pressure of 1MPa. The material is mild steel with a tensile strength of 500MPa. Assuming an efficiency of the longitudinal welded joint to be 75% and a factor of safety of 5 find the stress in the perforated steel plate.

A.1: Considering that the boiler design is based on thin cylinder principles the shell thickness is given by

\[ t = \frac{pr}{\sigma_y \eta} \]

where \( r \) is the boiler radius and \( \eta \) is the joint efficiency.

This gives

\[ t = \frac{10^6 \times 1.25}{\left(\frac{500}{5}\right) \times 10^6 \times 0.75} = 0.0166 \text{m} = 16.6 \text{mm}, \text{say } 20 \text{mm}. \]

The stress in the perforated plate is therefore given by \( \sigma = \frac{pr}{t} \) i.e. 62.5MPa

Q.2: A hydraulic cylinder with an internal diameter 250mm is subjected to an internal pressure of 10 MPa. Determine the wall thickness based on (a) Maximum principal stress theory, b) Maximum shear stress theory and c)
Maximum distortion energy theory of failure. Compare the results with wall thickness calculated based on thin cylinder assumption. Assume the yield stress of the cylinder material to be 60 MPa.

A.2:

Considering that the hydraulic cylinders are normally designed on the thick cylinder assumption we have from section 9.3.1.1 for Maximum Principal stress Theory we have

$$t = r_i \left( \frac{1}{1 + \frac{p_i}{\sigma_{yp}}} - 1 \right)$$

Here $$\frac{p_i}{\sigma_{yp}} = 10/60 = 0.167$$ and $$r_i = 125$$ mm. This gives $$t = 22.9$$ mm, say 23 mm.

From section 9.3.1.2 for Maximum Shear Stress theory we have

$$t = r_i \left( \frac{1}{1 - 2 \left( \frac{p_i}{\sigma_{yp}} \right)} - 1 \right)$$

With $$\frac{p_i}{\sigma_{yp}} \approx 0.167$$ and $$r_i = 125$$ mm, $$t = 28.2$$ mm, say 29 mm.

From section 9.3.1.4 for maximum distortion energy theory we have

$$t = r_i \left( \frac{1}{1 - \sqrt{3} \left( \frac{p_i}{\sigma_{yp}} \right)} - 1 \right)$$

with $$\frac{p_i}{\sigma_{yp}} \approx 0.167$$ and $$r_i = 125$$ mm $$t = 23.3$$ mm, say 24 mm.
Considering a thin cylinder $t = r_i \left( \frac{p_i}{\sigma_{yp}} \right)$ and this gives $t = 20.875\text{mm}$, say $21\text{mm}$.

The thin cylinder approach yields the lowest wall thickness and this is probably not safe. The largest wall thickness of $29\text{mm}$ predicted using the maximum shear stress theory is therefore adopted.

**Q.3:** A cylinder with external diameter $300\text{mm}$ and internal diameter $200\text{mm}$ is subjected to an internal pressure of $25\text{ MPa}$. Compare the relative merits of a single thick walled cylinder and a composite cylinder with the inner cylinder whose internal and external diameters are $200\text{mm}$ and $250\text{ mm}$ respectively. A tube of $250\text{ mm}$ internal diameter and $300\text{mm}$ external diameter is shrunk on the main cylinder. The safe tensile yield stress of the material is $110\text{ MPa}$ and the stress set up at the junction due to shrinkage should not exceed $10\text{ MPa}$.

**A.3:**

We first consider the stresses set up in a single cylinder and then in a composite cylinder.

**Single cylinder**

The boundary conditions are

at $r = 150\text{mm} \sigma_r = 0$ and at $r = 100\text{mm} \sigma_r = -20\text{MPa}$

Using equation (10) in section 9.2.1

$$C_1 + \frac{C_2}{0.0225} = 0 \quad \text{and} \quad C_1 + \frac{C_2}{0.01} = -20$$

This gives $C_1 = 16$ and $C_2 = -0.36$

The hoop stress at $r = 100\text{mm}$ and $r = 150\text{ mm}$ are $52\text{ MPa}$ and $32\text{ MPa}$ respectively.
**Stress in the composite cylinder**

The stresses in the cylinder due to shrinkage only can be found using the following boundary conditions

at \( r = 150\text{mm} \) \( \sigma_r = 0 \) and \( r = 125\text{mm} \) \( \sigma_r = -10\text{MPa} \)

Following the above procedure the hoop stress at \( r = 150 \text{ mm} \) and \( r = 125\text{mm} \) are **45.7 MPa** and **55.75 MPa** respectively.

The stress in the **inner cylinder** due to shrinkage only can be found using the following boundary conditions

at \( r = 100\text{mm} \) \( \sigma_r = 0 \) and \( r = 125\text{mm} \) \( \sigma_r = -10\text{MPa} \)

This gives the hoop stress at \( r = 100\text{mm} \) and \( r = 125\text{mm} \) to be **-55.55 MPa** and **-45.55 MPa** respectively.

Considering the internal pressure only on the complete cylinder the boundary conditions are

at \( r = 150\text{mm} \) \( \sigma_r = 0 \) and \( r = 100\text{mm} \) \( \sigma_r = -25 \text{MPa} \)

This gives

\[
\begin{align*}
(\sigma_\theta)_{r=150\text{mm}} &= 40 \text{MPa} \\
(\sigma_\theta)_{r=125\text{mm}} &= 49 \text{MPa} \\
(\sigma_\theta)_{r=100\text{mm}} &= 65 \text{MPa}.
\end{align*}
\]

Resultant stress due to both shrinkage and internal pressure

**Outer cylinder**

\[
(\sigma_\theta)_{r=150\text{mm}} = 40 + 45.7 = 85.7 \text{ MPa}
\]

\[
(\sigma_\theta)_{r=125\text{mm}} = 49 + 55.75 = 104.75 \text{ MPa}
\]

**Inner cylinder**

\[
(\sigma_\theta)_{r=125\text{mm}} = 49 - 45.7 = 3.3 \text{ MPa}
\]

\[
(\sigma_\theta)_{r=100\text{mm}} = 65 - 55.75 = 9.25 \text{ MPa}
\]

The stresses in both the single cylinder and the composite are within the safe tensile strength of the material. However in the single cylinder the stress gradient is large across the wall thickness whereas in the composite cylinder the stress variation is gentle. These results are illustrated in **figure- 9.3.2.1**
9.3.2.1F - Stress gradients (circumferential) in the inner and outer cylinders as well as the gradient across the wall of a single cylinder.

9.3.3 Summary of this Lesson

The lesson initially discusses the application of different failure theories in thick walled pressure vessels. Failure criterion in terms of the ratio of wall thickness to the internal radius and the ratio of internal pressure to yield stress have been derived for different failure criterion. Failure criterion for prestressed composite cylinders and cylinders with autofrettage have also been derived. Finally comparisons of different failure criterion have been discussed.

9.3.4 References for Module-9


Source: