

## Maxwell's Equations

We have considered Maxwell's curl equations for electrostatic fields and modified for time-varying situations to satisfy Faraday's Law. i.e  $\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$ .

We shall now consider Maxwell's curl equation for magnetic fields (Ampere's Law) for time-varying conditions.

For static fields EM fields, we have

$$\nabla \times \bar{H} = \bar{J}$$

$\bar{J}$ , includes conduction and convection currents.

$$\bar{J} = \sigma \bar{E} + \rho \bar{v}$$

$\sigma \bar{E}$  = Conduction current, due to the presence of the electric field in the conducting medium,

$\rho \bar{v}$  = Convection current, due to the motion of a free charge distribution.

The equations are:

$$\begin{aligned}\nabla \times \bar{E} &= -\frac{\partial \bar{B}}{\partial t} & \nabla \times \bar{H} &= \bar{J} \\ \nabla \cdot \bar{D} &= \rho & \nabla \cdot \bar{B} &= 0\end{aligned}$$

Although the four Maxwell's equations are consistent, they are not independent. The two divergence equations can be derived from the two curl equations by making use of the Continuity Equation  $\nabla \cdot \bar{J} = -\frac{\partial \rho}{\partial t}$  and Continuity Equation can be derived from four Maxwell's Equations.

Consider  $\nabla \times \bar{H} = \bar{J}$ , and  $\nabla \cdot \nabla \times \bar{H} = \nabla \cdot \bar{J} = 0$ ,

but the continuity equation is,

$$\nabla \cdot \bar{J} = -\frac{\partial \rho}{\partial t}$$

for time-varying fields. Ampere's Law equation needs modification. Add a term to the Ampere's Law equation:

$$\nabla \times \bar{H} = \bar{J} + \bar{J}_d$$

and operate  $\nabla \cdot$  on both sides:

$$\nabla \cdot (\nabla \times \bar{H}) = \nabla \cdot \bar{J} + \nabla \cdot \bar{J}_d$$

$$0 = -\frac{\partial \rho}{\partial t} + \nabla \cdot \bar{J}_d$$

Since,

$$\nabla \cdot \bar{D} = \rho$$

$$\nabla \cdot \left( \bar{J}_d - \frac{\partial \bar{D}}{\partial t} \right) = 0$$

Let,  $\bar{J}_d = \frac{\partial \bar{D}}{\partial t}$ , Displacement current density, A/m<sup>2</sup>. The Ampere's Law equation for time-varying fields takes the form:

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}$$

Then four consistent equations,

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} \quad \text{Faraday Law}$$

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t} \quad \text{Ampere's Law}$$

$$\nabla \cdot \bar{D} = \rho \quad \text{Gauss Law}$$

$$\nabla \cdot \bar{B} = 0 \quad \text{Magnetic Flux Law}$$

are known as Maxwell's equations.

Maxwell's equations, together with

$$\bar{F} = q(\bar{E} + \bar{u} \times \bar{B}) \quad \text{Lorentz Force Equation}$$

$$\nabla \cdot \bar{J} + \frac{\partial \rho}{\partial t} = 0 \quad \text{Continuity Equation}$$

from the foundation of the electromagnetic theory.

## INTEGRAL FORMS OF MAXWELL'S EQUATIONS

Consider  $\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$  and integrate over an open surface  $S$  with a contour  $C$  and apply Stoke's Theorem:

$$\int_S \nabla \times \bar{E} \cdot \hat{n} ds = - \int_S \frac{\partial \bar{B}}{\partial t} \cdot \hat{n} ds$$

Now consider,

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}$$

And integrate over a surface  $S$

$$\int_S \nabla \times \bar{H} \cdot \hat{n} ds = \int_S \left( \bar{J} + \frac{\partial \bar{D}}{\partial t} \right) \cdot \hat{n} ds$$

$$\oint_C \bar{H} \cdot d\bar{l} = \int_S \left( \bar{J} + \frac{\partial \bar{D}}{\partial t} \right) \cdot \hat{n} ds$$

$$\oint_C \bar{H} \cdot d\bar{l} = I + \int_S \left( \frac{\partial \bar{D}}{\partial t} \right) \cdot \hat{n} ds$$

For the divergence equations:

Consider  $\nabla \cdot \bar{B} = 0$ ,

Take the volume integral of both sides of the divergence equation over a volume  $V$  and apply the divergence theorem:

$$\int_V (\nabla \cdot \bar{B}) dv = 0$$

results,

$$\oint_S \bar{B} \cdot \hat{n} ds = 0$$

And similarly,  $\nabla \cdot \bar{D} = \rho$

$$\int_V (\nabla \cdot \bar{D}) dv = \int_V \rho dv$$

$$\oint_S \bar{D} \cdot d\bar{s} = Q$$

Consequently, the boundary conditions remain valid for the time-varying fields, where  $\hat{a}_n$  is the unit vector normal to the boundary.

$$\begin{aligned} E_{2_t} &= E_{1_t} & \text{or} & & (\bar{E}_1 - \bar{E}_2) \times \hat{a}_n &= 0 \\ H_{1_t} - H_{2_t} &= J & \text{or} & & (\bar{H}_1 - \bar{H}_2) \times \hat{a}_n &= \bar{J} \\ D_{1_n} - D_{2_n} &= \rho_s & \text{or} & & (D_1 - D_2) \cdot \hat{a}_n &= \rho_s \\ B_{1_n} - B_{2_n} &= 0 & \text{or} & & (B_1 - B_2) \cdot \hat{a}_n &= 0 \end{aligned}$$

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